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Important physical properties of solids are directly related to the capacity of these materials to undergo phase transitions, resulting in change of crystal morphology, during their manufacture or use. New properties are attainable in materials by exploiting this morphology change. Insight into the physical basis of phase transition processes is available through Landau theory.

Several important materials such as zeolites, perovskites, memory metals and superconductors undergo phase transitions. Zeolite ZSM-5 has been shown to undergo a phase transition from a monoclinic structure to an orthorhombic structure at approximately 40°C [1,2]. Because this material is grown at elevated temperatures where it exists as the orthorhombic phase, severe twinning occurs when it is cooled to the monoclinic phase at room temperature. This twinning makes a detailed structural study of ZSM-5 difficult. The perovskite BaTiO₃, an important ferroelectric material, undergoes a thermally induced tetragonal to cubic transition at 125°C accompanied by a transition from ferroelectric to paraelectric [3]. TiNi, a material used to join titanium hydraulic tubing in the Grumman F14 Navy jet fighters, undergoes a monoclinic to cubic transition which is directly linked to its most interesting property, that of shape memory [5]. The property of superconductivity is closely related to the structure of the material. $Zr_{1-x}S$ is a low-temperature superconductor for which the low-temperature structure was not known until fairly recently [11,12,17]. From the aforementioned examples it is clear that the structure and properties of materials are directly linked.

A phase transition of one material may be used to engineer different properties into other materials such as is done in transformation toughening. Here, ZrO_2 is introduced into an AI₂O₃ matrix, a brittle material in its pure form, to make the ceramic more shock-resistent [4]. In this application the ZrO_2 acts as a toughener by absorbing stresses incurred upon the matrix by undergoing a stress-induced phase transition.

Due to the important and often difficult task of understanding phase transitions, it is fortunate that a strong theoretical framework exists, the Landau theory [6-9] of symmetry and phase transitions. It is not uncommon for it to be very difficult, if not impossible, to unambiguously determine the structure of a solid which has undergone a phase transition due to resultant twinning and the low intensity of diffraction lines arising from superstructure. The use of Landau theory to clarify diffraction patterns has been very recently reviewed by H. F. Franzen [10].

The symmetry rules of Landau theory only apply to second-order phase transitions, which are necessarily continuous. First- and second-order transitions vary in that a transition in which there does not exist a region where both phases coexist in equilibrium is necessarily second-order. This property suggests that these transitions are strongly bound by symmetry. The main thrust of Landau's theory is that it is possible to expand the Gibbs free energy function in a Taylor series as a function of the order parameter, η , where η describes the overall extent of the distortion (in the case of an order-disorder transition, $\eta = 0$ implies a perfectly disordered structure and $\eta = 1$ describes a completely ordered structure).

$$G = G^{\circ} + \alpha \eta + A \eta^2 + B \eta^3 + C \eta^4 + ...$$

 α must be zero at $\eta = 0$ for G to be a minimum at this point

$$G = G^\circ + A\eta^2 + B\eta^3 + C\eta^4$$

where A, B and C are functions of state variables. If B does not equal 0 by symmetry it is impossible for a second-order transition to occur except at isolated state points where B goes to zero by chance. Therefore, the Gibbs free energy function which describes a second-order transition is:

$$G = G^{\circ} + A\eta^2 + C\eta^4$$

where C has been taken to be positive (if C were less than zero the function would have to be expanded to the sixth order which necessarily leads to a first-order transition). A second-order transition occurs as A passes through zero at the transition point. It follows that for A to go to zero at a succession of points, A must correspond to a single irreducible representation. Because G must be the same for equivalent distortions it must be invarient under symmetry operations. Therefore, this function can be written as a function of invarient combinations of the basis functions which correspond to the transition. The basis functions which describe the transition are those which are symmetric to the symmetry operations retained and anti-symmetric to those which are lost. This function may be minimized to give possible space groups which can result from a second-order transition from some parent symmetry.

This approach along with experimental observations has been used in the examination of the defect NaCl structure of $Zr_{1-x}S$ [11,12], $Sc_{1-x}S$ [11], the oxides of Ni, Mn, Co and Fe [10,13-16] of which $Zr_{1-x}S$, $Sc_{1-x}S$, and others are low-temperature superconductors [17]. This technique has been used in numerous intermetallics with the CsCl parent structure such as RhTi [18,19], NbRu_{1+x} [19], TaRu [20], AuCd [21], TiNi [21], V_{0.54}Ir_{0.46} [23] and LiPb [21], several of which are related to memory metals. J. D. Corbett has recently used the theory to help explain the formation of unequal Zr-Zr bonds in ZrI₃ [22,24].

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